

Efficient Exploration of Membrane-Associated Phenomena at Atomic Resolution

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Abstract Biological membranes constitute a critical component in all living cells. In addition to providing a conducive environment to a wide range of cellular processes, including transport and signaling, mounting evidence has established active participation of specific lipids in modulating membrane protein function through various mechanisms. Understanding lipid–protein interactions underlying these mechanisms at a sufficiently high resolution has proven extremely challenging, partly due to the semi-fluid nature of the membrane. In order to address this challenge computationally, multiple methods have been developed, including an alternative membrane representation termed highly mobile membrane mimetic (HMMM) in which lateral lipid diffusion has been significantly enhanced without compromising atomic details. The model allows for efficient sampling of lipid–protein interactions at atomic resolution, thereby significantly enhancing the effectiveness of molecular dynamics simulations in capturing membrane-associated phenomena. In this review, after providing an overview of HMMM model development, we will describe briefly successful application of the model to study a variety of membrane processes, including lipid-dependent binding and insertion of peripheral proteins, the mechanism of phospholipid insertion into lipid bilayers, and characterization of optimal tilt angle of transmembrane helices. We conclude with practical recommendations for

proper usage of the model in simulation studies of membrane processes.

Keywords Cellular membrane · Lipid bilayer · Membrane proteins · Peripheral proteins · Molecular dynamics simulation

Introduction

Biological membranes are an essential component of all living cells (Gennis 1989; Mouritsen 2005). The cellular membrane is a complex and dynamic environment with a large diversity of lipids that shape its characteristics (van Meer et al. 2008; van Meer and de Kroon 2011; Ingólfsson et al. 2014). It has been estimated that several hundred distinct types of lipids populate the typical mammalian plasma membrane (Sampaio et al. 2011; Klose et al. 2013), influencing membrane fluidity, geometry, pressure, and surface charge (Klose et al. 2013).

Approximately, 26 % of the human proteome is known to be functional only within membranes (Fagerberg et al. 2010) and these membrane proteins account for an estimated 60 % of current drug targets (Yildirim et al. 2007). Membranes, and the proteins embedded in the membrane, are essential to life. They actively establish and maintain electrochemical gradients, and further exploit these gradients to perform meaningful transport of nutrients and waste into and out of the cell. Membrane proteins are also indispensable to cellular signaling, as they provide the coupling mechanism between the inside and outside of the cell (Grecco et al. 2011). Peripheral membrane proteins, which contain structurally varying lipid-binding globular domains that recognize specific membrane components (Lemmon 2008), associate with the membrane surface and

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are involved in signaling, trafficking, and regulating cell structure.

The membrane itself can also be used as a signal. By modulating the lipid composition through leaflet dependent (van Meer et al. 2008; van Meer and de Kroon 2011) exposure of anionic phospholipids [e.g., phosphatidylserine (PS)] or biosynthesis of particular rare lipid species, such as phosphatidylinositides or cardiolipin, the activity of peripheral (Lemmon 2008) and transmembrane (TM) (Andersen and Keoppe 2007) proteins can be drastically altered. Thus, the importance of the membrane and its dynamic character in cellular homeostasis cannot be overstated.

Despite the importance of the membrane and membrane proteins to the maintenance and survival of the cell, the inherent fluid nature of the lipid bilayer (Singer and Nicolson 1972; Engelman 2005) makes detailed experimental studies on the specific roles of lipids in membrane proteins extremely challenging. This problem is compounded for peripheral membrane proteins. Due to the reversible nature of the interaction between peripheral proteins and the membrane, crystallizing the membrane-bound structure of peripheral proteins is exceedingly difficult. While techniques such as SAXS (Denisov et al. 2004; Mosbaek et al. 2010; Elliott et al. 2010), EPR (Lin et al. 1998; Kohout et al. 2003), NMR (Goult et al. 2009), FRET (McCallum et al. 1997; Kohout et al. 2003), and mutagenesis studies can help determine the approximate binding face of peripheral proteins, more structural information is needed to reveal molecular details critical to membrane binding of peripheral proteins and the role of lipids within this binding.

Molecular dynamics (MD) simulations offer sufficient spatial and temporal resolutions to capture specific lipid–protein interactions and have been extensively used for this purpose (Ohkubo and Tajkhorshid 2008; Lindahl and Sansom 2008; Ayton and Voth 2009; Lai et al. 2010; Khalid and Bond 2013; Bucher et al. 2013; Braun et al. 2014). Due to the slow dynamics of lipids (Klauda et al. 2006; Wohllert and Edholm 2006) on time scales accessible to conventional atomistic MD simulations, the membrane is effectively static, preventing one from sufficiently describing the lipid motion necessary for lipid mixing and adequate sampling of lipid–protein interactions. Alternative computational approaches have been devised to circumvent the problem of slow lipid dynamics, including coarse-grained (Shelley et al. 2001; Marrink et al. 2004; Izvekov and Voth 2005; Marrink et al. 2007; Ayton et al. 2010; Marrink and Tieleman 2013) and implicit (Feig and Brooks 2004; Im and Brooks 2005; Chen et al. 2006; Bu et al. 2007; Mondal et al. 2010) membrane models. These approaches, which sacrifice atomistic detail for accelerated dynamics, are of limited utility in resolving lipid-specific interactions vital to membrane protein function.

In order to circumvent the aforementioned technical issues precipitated by slow lipid diffusion, a recently developed atomic membrane model, termed the highly mobile membrane mimetic (HMMM) (Arcario et al. 2011; Ohkubo et al. 2012), provides a suitable alternative for studying protein–lipid interactions. The HMMM model first fully developed by Ohkubo et al. (2012) is constructed using a combination of surfactant-like short-tailed lipids, representing *in full atomic detail* the lipid head groups that are often the main protein-interacting elements, with a liquid representation of the membrane core (Fig. 1). The HMMM representation significantly accelerates lateral lipid diffusion and enhances lipid–protein sampling (Arcario et al. 2011; Ohkubo et al. 2012). A particular strength of this membrane representation is that it affords multiple simulations of spontaneous protein interactions with the membrane in less time than is required to simulate the biased binding of a peripheral protein in a conventional simulation (Ohkubo et al. 2012; Baylon et al. 2013; Vermaas and Tajkhorshid 2014a; Arcario and Tajkhorshid 2014), providing for efficient and enhanced sampling of protein–lipid interactions. By design, the kinetics of membrane-associated processes are accelerated by the HMMM model due to faster lipid diffusion. As far as thermodynamical properties are concerned, the energetics associated with protein–membrane interactions is adequately captured for the interfacial and surface regions

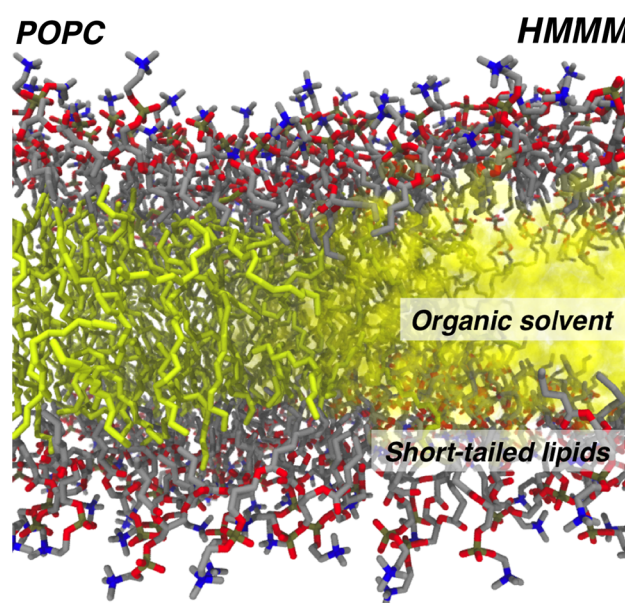


Fig. 1 Comparison of a full-tail and HMMM membrane representations pictorially illustrated by gradual transformation of a POPC membrane (left) into an HMMM membrane (right). The hydrophobic core of the membrane (yellow) is replaced by an organic solvent in the HMMM membrane, while the headgroup and short tails of the lipids (typically first five carbons) are preserved (Color figure online)

of lipid bilayers, but significant deviations are evident in the core of the membrane due to the fluid nature and polarity of the solvent used to replace the bulk of lipid tails (Pogorelov et al. 2014).

The HMMM model has been successfully applied to simulation studies of several membrane-associated protein systems (Ohkubo et al. 2012; Baylon et al. 2013; Vermaas and Tajkhorshid 2014a; Blanchard et al. 2014; Arcario and Tajkhorshid 2014; Wu and Schulten 2014; Rhéault et al. 2015). We begin this review by describing the development trajectory of the HMMM model, followed by a set of its recent applications to a wide variety of membrane-associated phenomena including: phospholipid insertion into membrane (Vermaas and Tajkhorshid 2014b), binding and insertion of peripheral proteins such as cytochrome P450 (Baylon et al. 2013), hemoglobin N (Rhéault et al. 2015), talin (Arcario and Tajkhorshid 2014), synaptotagmin I (Syt1) (Wu and Schulten 2014), and α -synuclein (Vermaas and Tajkhorshid 2014a), and studies of pre-fusion configuration of synaptobrevin TM helix in a lipid bilayer (Blanchard et al. 2014). We conclude with outlining the best practices for utilizing the HMMM model and an outlook on its future developments.

Development Trajectory of the HMMM

The initial development of the HMMM was driven by a need to study membrane-binding processes and other phenomena at the membrane surface that were inaccessible to MD simulations. The test case that accompanied the initial development of the HMMM is the GLA domain, the membrane-binding domain of the vitamin K-dependent clotting factors. The GLA domain is named for its characteristic vitamin K-dependent γ -carboxyglutamate (Gla) residues that tightly coordinate multiple Ca^{2+} ions. The coordinated Ca^{2+} ions then serve two functions: they play a structural role, maintaining the “active” form of the GLA domain, as well as an adhesive role, interacting with anionic phospholipid headgroups, including PS (Ohkubo and Tajkhorshid 2008). Despite the established importance of PS lipids in the process, it was unclear which moieties of the GLA domain interacted with these lipids and to what degree the protein is inserted in the membrane; while Nelsestuen (1999) proposed the N-terminal ω -loop and K32 of the GLA domain made interfacial contact with the membrane, Falls et al. (2001) and Mizuno et al. (2001) suggested that only residues in the hydrophobic “keel” were inserted into the membrane. Given the overly simplistic view of the membrane going into these postulated models, they turned out to only partially capture the

interactions involved in membrane binding of the GLA domain.

Ohkubo and Tajkhorshid (2008) explored membrane binding of the fVII GLA using conventional atomistic MD simulations with two alternative methods of triggering membrane binding, and incorporated available experimental data (McCallum et al. 1996; Waters et al. 2006; Colina et al. 2006) into the modeling process. The first method was to remove lipids from the membrane patch below the inserting species, and through successive minimization stages, forcibly insert the protein domain into the bilayer. In the second, the GLA domain was inserted into the bilayer through the use of steered MD (SMD) (Grubmüller et al. 1996; Leech et al. 1996; Izrailev et al. 1997), in which an external applied force drives the insertion over the course of the simulation. The resulting model was unique, in that both the hydrophobic ω -loop and Ca^{2+} ions contributed substantial interactions with the membrane.

While the devised approach successfully produced the most detailed model of a membrane-bound GLA domain, this approach was not broadly applicable to studying the insertion of other proteins that interact with the membrane. Without a target depth, SMD can easily result in a biased membrane-bound configuration and, due to the slow relaxation time of lipids, equilibration of several hundred nanoseconds or longer may be required to unbiased the simulation system. Furthermore, without an applied bias, Ohkubo and Tajkhorshid (2008) did not observe GLA domain insertion into the membrane in over 350 ns of simulation. Finally, the model also depended on a homogeneous lipid composition, a pure PS bilayer, as again, lateral lipid exchange is a slow process on atomistic MD timescales, and the mixing time for a realistic membrane is rather large, estimated to be on the microsecond timescale by Ingólfsson et al. (2014) through coarse-grained simulations.

In order to address these sampling shortcomings, Arcario et al. (2011) tested atomistic biphasic solvent systems for their ability to facilitate rapid partitioning of the human protein C (hPrC) GLA domain, an anticoagulant, at the interface between water and a hydrophobic phase. In these biphasic models, hydrophobic compounds whose parameters were available in the CHARMM general force field (Vanommeslaeghe et al. 2010) were constructed into a layer designed to represent the hydrophobic membrane core. The optimal compound had to be largely hydrophobic, liquid under typical simulation conditions, and preferably have a large diffusion constant to enable the membrane constituents to quickly reorganize around an inserting species. To a first approximation, the diffusion constant is inversely proportional to molecular size

(Einstein 1905); however, small hydrophobic compounds are gaseous. This poses a difficult optimization problem, as no natural compound is simultaneously small, hydrophobic, and liquid at ambient conditions, and so a compromise needs to be struck between these three competing goals (Fig. 2). Arcario et al. (2011) tested four compounds, including 1,1-dichloroethane (DCLE), dimethyl sulfide (DMS), ethyl propyl ether (EPE), and heptane, which all represent approximate solutions to this problem.

Partitioning of the GLA domain to the water/DCLE interface and partial insertion into the organic phase was fast, occurring spontaneously in under 10 ns, whereas for the other solvents tested, complete and permanent insertion was not observed (Arcario et al. 2011). Additionally, DCLE showed the largest diffusion constant among the compounds tested (Fig. 2), making it the clear choice for use as the membrane core mimic in biphasic simulation system studies. This biphasic insertion model, where a bilayer is modeled as a box of organic solvent, formed the initial step towards the development of the HMMM model.

What was added into the biphasic model formulation by Ohkubo et al. (2012) are two layers of surfactants (short-tailed lipids), whose headgroups are identical to those found in physiological phospholipids, but where the acyl tails are truncated beyond the fifth carbon atom. The surfactants spontaneously partition to the interface of the DCLE-water biphasic system, generating a membrane mimetic (Ohkubo et al. 2012). When used to study the insertion of GLA domains, the mimetic enables not only the rapid spontaneous binding of the GLA domain to the membrane, but also recovers the specific Ca^{2+} -headgroup interactions (Fig. 3). This result, combined with the favorable enhanced lipid lateral diffusion and headgroup

atomic distributions (Ohkubo et al. 2012), began a push in the field to apply this methodology to proteins which are known or are thought to have substantial membrane interaction (Baylon et al. 2013; Wu and Schulten 2014; Arcario and Tajkhorshid 2014; Blanchard et al. 2014; Vermaas and Tajkhorshid 2014a; Rhéault et al. 2015).

The potential to apply the HMMM model to other membrane-interacting peptides spurred further work, specifically in validating the protein-membrane interaction energetics of the model (Pogorelov et al. 2014). For determining the protein-membrane interaction characteristics, the procedure followed by Pogorelov et al. (2014) closely matches that done in prior atomistic (MacCallum et al. 2008) and coarse-grained (Monticelli et al. 2008) computational studies, where the free energy profile (i.e., the potential of mean force, abbreviated PMF) along the membrane normal was computed for amino acid side chain analogs using umbrella sampling (Torrie and Valleau 1977) and WHAM (Kumar et al. 1992). In general, the accuracy of the HMMM energetics relative to a conventional full membrane representation are on par with or superior to coarse-grained representations (Fig. 4). The accuracy is best on the membrane periphery, where the short-tailed lipids fully represent the atomic structure of a conventional bilayer. In the membrane interior, the fluid nature of the solvent comprising the hydrophobic core allows extra rotational degrees of freedom for bulkier side chains (e.g., Phe and Trp; Fig. 4b) to be populated, which stabilizes them in these regions beyond what is observed in a full lipid bilayer. Additionally, the DCLE solvent has a larger dipole than the lipid acyl tails it is replacing, and as a result stabilizes polar or charged side chains in the membrane interior to a greater extent than is observed in a

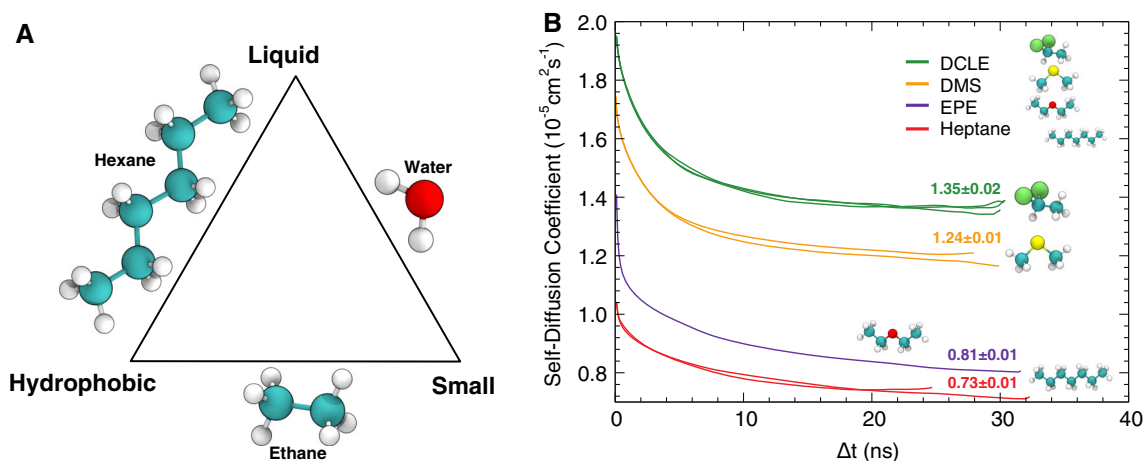
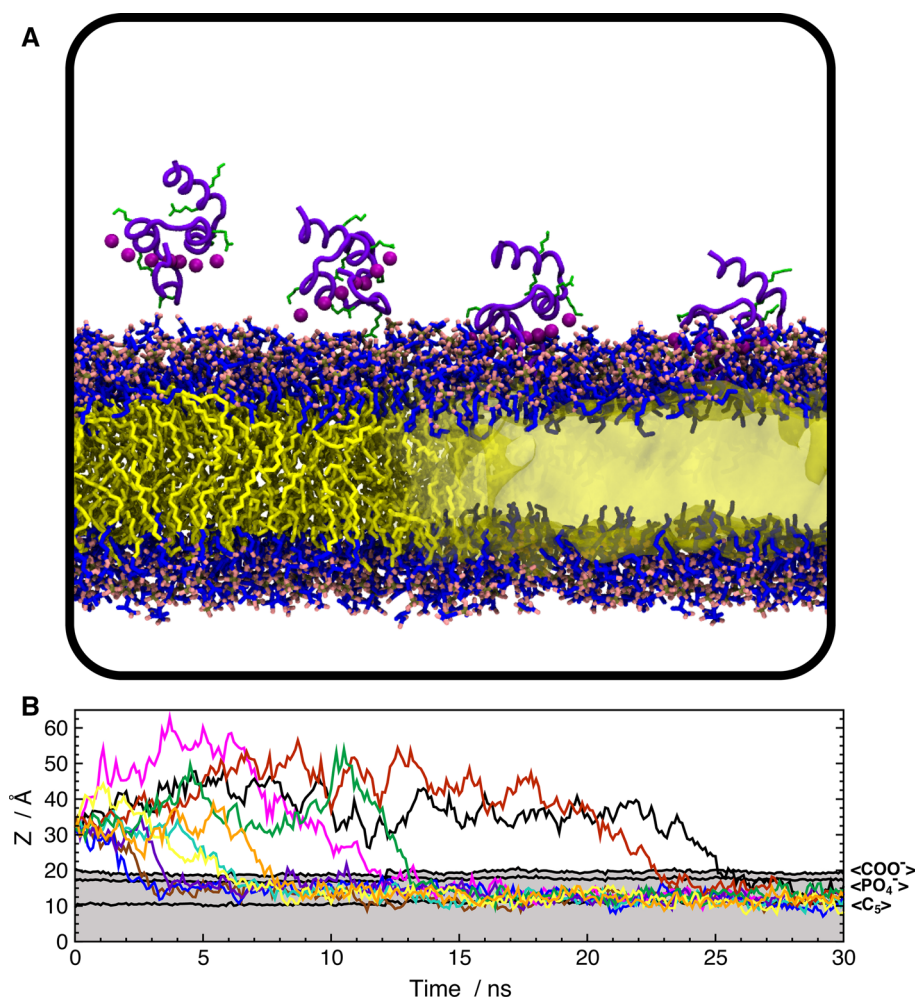


Fig. 2 **a** The optimization triangle for desirable membrane mimics, emphasizing that natural compounds can only satisfy two of the three desired properties. **b** Self-diffusion constants for tested membrane mimics. Smaller compounds diffuse more rapidly, thereby quickly

accommodating protein insertion. DCLE is 1,1-dichloroethane, DMS is dimethyl sulfide, and EPE is ethyl propyl ether. Adapted from Arcario et al. (2011)

Fig. 3 Spontaneous binding and insertion of the GLA domain to anionic membranes captured by HMMM. **a** The binding of GLA domains (purple trace) to PS headgroups is mediated by bound Ca^{2+} ions (purple spheres) and basic side chains (green licorice). Ohkubo and Tajkhorshid (2008) and Ohkubo et al. (2012) studied the binding process in both conventional (left) and HMMM (right) bilayers. **b** Penetration depth over time for 10 independent HMMM binding simulations, in which all GLA domains bind to the target membrane within 30 ns and converged to a single insertion depth (Color figure online)



conventional bilayer. These solvent artifacts can be responsibly managed by geometric constraints on the protein to prevent solvent intercalation where appropriate.

Applications of HMMM to Membrane-Associated Phenomena

What makes these artifacts tolerable is the degree to which the HMMM model accelerates membrane-associated phenomena. Vermaas and Tajkhorshid (2014b) studied a model insertion system, which demonstrated that HMMM bilayers accelerate insertion of a single phospholipid by a factor of 3 by taking advantage of their higher lateral lipid diffusion to “welcome” the inserting species. In membrane-interacting peptide studies where the inserting species is larger, the impact of the accelerated diffusion is enhanced by 1–2 orders of magnitude, depending on the degree of membrane interaction with the inserting species (Ohkubo et al. 2012; Vermaas and Tajkhorshid 2014a).

The accelerated lipid diffusion and resulting rapid insertion have been used to study membrane binding and membrane activity of various membrane-associated processes (Fig. 5). The applications include developing a model for phospholipid adsorption and desorption mechanics (Vermaas and Tajkhorshid 2014b), membrane-induced conformational changes that modulate drug and oxygen entry pathways of cytochrome P450 (Baylon et al. 2013) or hemoglobin (Rhéault et al. 2015), and drive domain rearrangement of the F2F3 domains of talin (Arcario and Tajkhorshid 2014). The HMMM has also been applied to several components of the cellular vesicle fusion machinery, including the Ca^{2+} sensitive synaptotagmin-1 (Wu and Schulten 2014), the SNARE component synaptobrevin (Blanchard et al. 2014), and the flexible amphipathic helix α -synuclein (Vermaas and Tajkhorshid 2014a). These applications will be examined in turn for their key findings in addition to the lessons that can be applied to other systems.

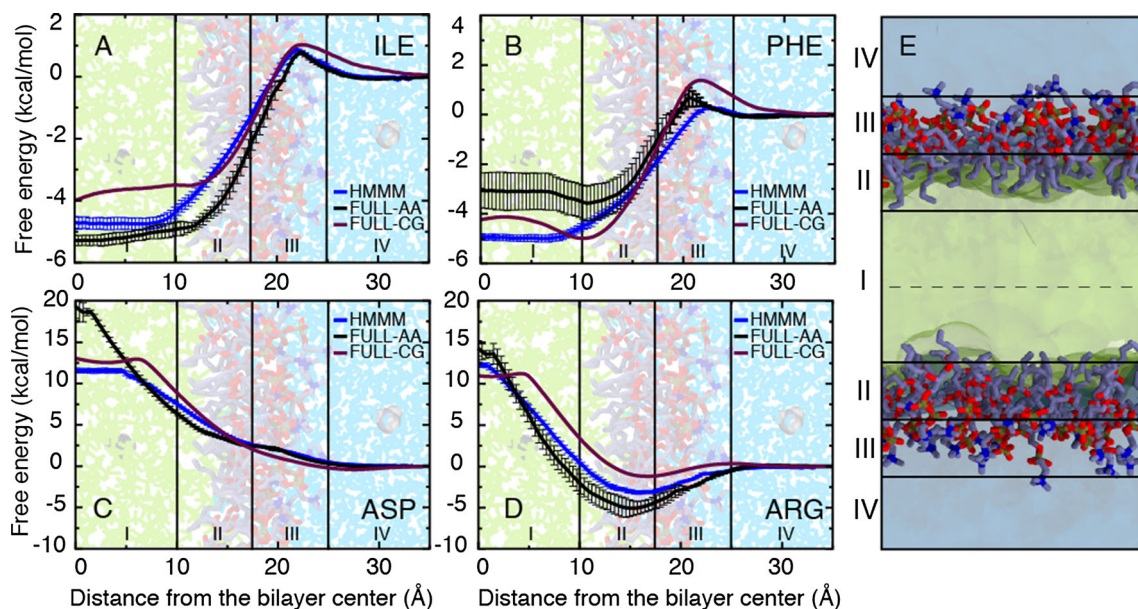
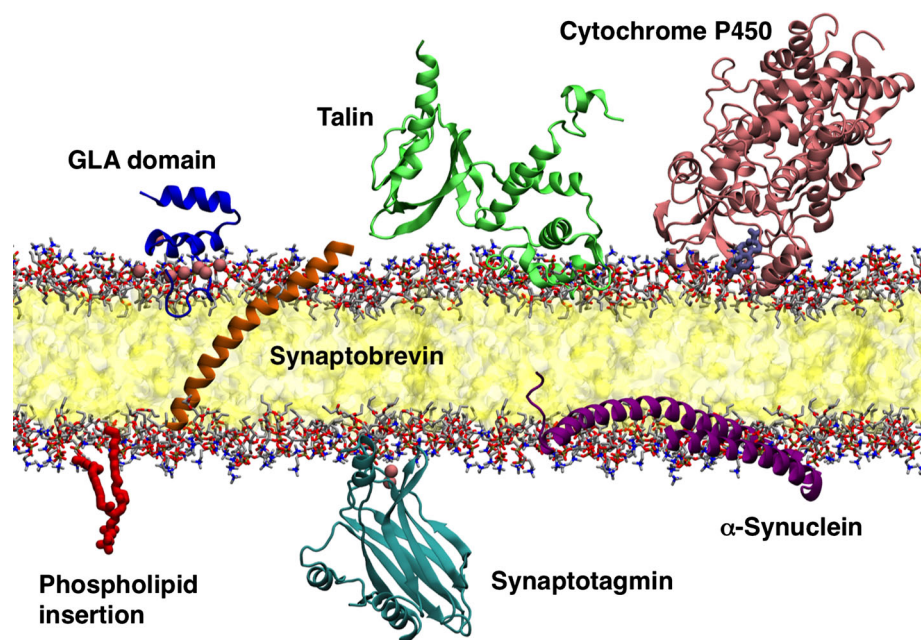


Fig. 4 PMFs for representative side chain analogs: aliphatic—Ile (a), hydrophobic—Phe (b), and charged—Asp (c) and Arg (d). Data are presented for the HMMM membrane (blue), FULL-AA (black) (MacCallum et al. 2008), and FULL-CG (purple) (Monticelli et al.

2008). All PMFs are presented by setting the free energy to zero in aqueous solution. Membrane regions I–IV are defined in the panel (e). Error estimates were obtained using bootstrap analysis. Data from Pogorelov et al. (2014) (Color figure online)

Fig. 5 Pictorial summary of systems reviewed here studied using the HMMM membrane



Phospholipid Insertion Mechanics: “One Tail at a Time”

Phospholipid insertion studies were initially carried out to demonstrate the efficiency of the HMMM on a small test system where direct comparisons to conventional bilayer simulations are possible, the model for phospholipid insertion into biological membranes developed by Vermaas and

Tajkhorshid (2014b) also led to novel insight as to the mechanism of phospholipid adsorption from solution, the primary mechanism for phospholipid exchange at low concentrations (McLean and Phillips 1981; Nichols and Pagano 1981). This study was complementary to work by Tieleman and Marrink (2006) and Grafmüller et al. (2013), which used umbrella sampling simulations to determine the free energy of phospholipid unbinding using the position of the center of

mass of the adsorbing or desorbing lipid as the reaction coordinate. Through repeated binding of phospholipids to both conventional and HMMM bilayers, Vermaas and Tajkhorshid (2014b) developed a kinetic model, whereby it is demonstrated that realistic binding and unbinding pathways pass through distinct intermediate stages (Fig. 6a), consistent with splayed-tail intermediates observed by Smirnova et al. (2010) in vesicle fusion studies.

The stages enumerated by Vermaas and Tajkhorshid (2014b) are roughly analogous to the process of putting on pants in the morning, in that it occurs one acyl tail at a time. The phospholipid first associates to the bilayer, and may bind and unbind reversibly looking for a suitable insertion location, much like when the reader may sort through their pants to select their favorite pair. After association, the phospholipid inserts its acyl tails in sequence irreversibly in the same manner as one would step into a pair of pants. Finally, the phospholipid puts on its belt when the headgroup reorients and relaxes within the bilayer.

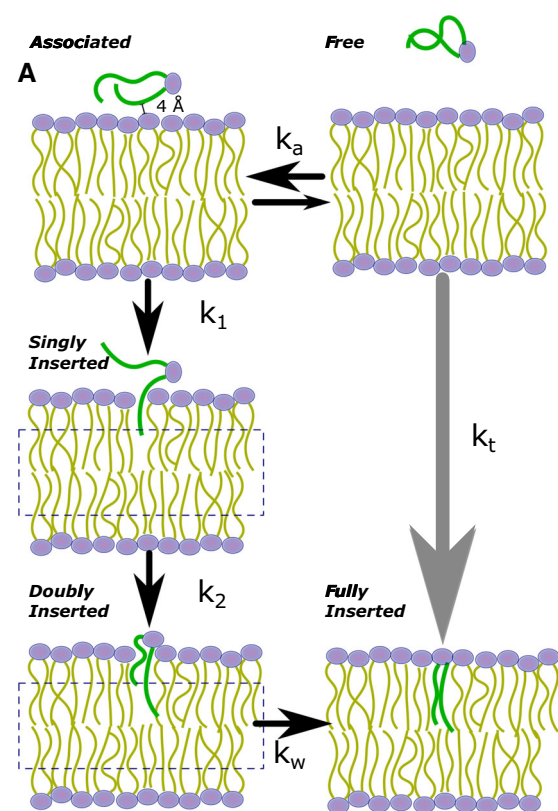
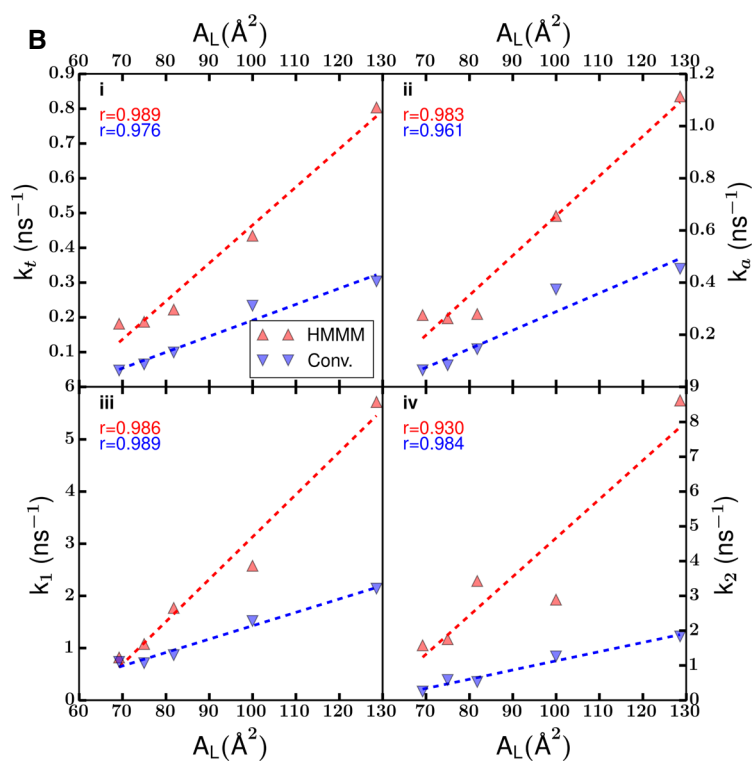


Fig. 6 **a** Schematic of the phospholipid insertion process, where the unbound phospholipid (green tails) associates with the membrane reversibly with rate k_a , thereafter inserts a single tail with rate k_1 , followed shortly thereafter by the insertion of the second tail at rate k_2 and headgroup reorganization at rate k_w . The overall process is governed by the summation of these steps, and occurs at an overall rate k_t . **b** Rate comparisons for individual steps from the schematic in

This model is consistent between both a conventional and a HMMM representation, however, by breaking down the process into individual steps, Vermaas and Tajkhorshid (2014b) also determined specifically what part of the overall process was accelerated, and how changing the membrane density impacts kinetics. Both reduced membrane density by increasing the area per lipid (A_L) and the switch to a HMMM bilayer accelerated the rate of insertion (Fig. 6b). However, digging deeper into the data, what is remarkable is where the acceleration for the HMMM bilayer is taking place. At physiological lipid densities, the HMMM accelerates the association step the most, leaving the rate of the first lipid insertion after association unchanged. The analysis of the second tail insertion rate is more complicated, as there are cases where the first tail inserted while the second acyl tail had not yet associated to the bilayer, and so the insertion rate of the second tail does not match. Vermaas and Tajkhorshid (2014b) concluded that HMMM preferentially accelerates the mechanistically simple part of insertion, namely the searching phase where



(a) between inserting into an HMMM and a conventional full bilayer over different areas per lipid (A_L). Each subpanel i–iv compares the rates for a different step, with a consistent color scheme across panels (blue for conventional bilayers, and red for HMMM bilayers). Adapted from Vermaas and Tajkhorshid (2014b). Copyright 2014 American Chemical Society (Color figure online)

the inserting species is sampling the bilayer surface for a gap to insert into. When this insight is applied to other systems, it suggests that the mechanism and interactions between the membrane and the inserting species are faithfully represented, with the insertion acceleration solely a result of the increased fluidity of the HMMM bilayer interface.

Increasing Drug Accessibility to Cytochrome P450 Through Membrane Binding

Cytochrome P450 (CYP) 3A4 is the most abundant isoform of the CYP family of enzymes in the human body, and it is involved in the metabolism of both xenobiotics, e.g., drugs, and endogenous compounds such as steroidal hormones. The large range of clinically available drugs metabolized by CYP3A4, more than 50 % (Guengerich 1999), has made this enzyme the target of extensive studies. Importantly, accumulating evidence suggests that interaction of CYP3A4 with the membrane is crucial for the recruitment of its lipophilic substrates to the active site of the enzyme (Williams et al. 2003; Schleinkofer et al. 2005), and therefore membrane binding and protein–lipid interactions might play critical roles in the function of the enzyme. However, details of these interactions were still elusive due to the lack of structural information concerning the membrane-bound form of CYP3A4.

Taking advantage of the enhanced lipid dynamics of the HMMM model, Baylon et al. (2013) characterized the membrane-bound form of the CYP3A4 by performing multiple independent simulations in which spontaneous binding of the enzyme to the membrane was consistently captured. Five independent simulations in which the globular domain of CYP3A4 (i.e., without the TM helix) was placed initially away from a PC membrane with different initial orientation were performed. In each case CYP3A4 was observed to bind to the membrane within 10 ns of the simulation. These simulations revealed structural elements located in the globular domain of CYP3A4 that mediate lipid interactions, and contribute to the insertion of the protein in the membrane. The convergence of the membrane-bound model of CYP3A4 was assessed by calculating the orientation of the protein in the membrane, measured as the relative angle of the heme moiety with respect to the membrane normal (Fig. 7), with an average of $72.2 \pm 3.2^\circ$ (Baylon et al. 2013). The calculated average angle supports the convergence of the membrane-bound model, since the final tilt angles deviate only by a few degrees, even after starting from initial orientations in solution that differed by 40° . Additional simulations of the membrane-bound CYP3A4 in which the TM helix was added to the protein revealed that the orientation of the protein is mainly determined by the direct interaction of the

globular domain with the membrane and not by the presence of the helix. Crucially, the orientation of HMMM bound CYP3A4 is comparable to linear dichroism measurements performed in parallel to the simulations for nanodisc-bound CYP3A4 (Baylon et al. 2013), where the average tilt angle was found to be $59.7 \pm 4.1^\circ$, yielding excellent agreement between simulations and experiment.

From the simulations, Baylon et al. (2013) observed that the interaction of CYP3A4 with the membrane favors transient opening of small access pathways leading to the active site by inducing conformational changes of side chains located at the protein–membrane interface. Interestingly, some of these access tunnels are not observed in the solution simulation of the enzyme, i.e., in the absence of the membrane. Based on their location, these tunnels may function as access pathways for compounds in or out of the CYP3A4 active site, and highlight the importance of correctly capturing lipid–protein interactions and atomic detail in understanding the mechanism of this key enzyme in the human body. The resulting membrane-bound structures can be used in subsequent simulations to further probe different functions of CYP3A4 in the membrane with atomic detail, such as the recently characterized mechanism of drug–drug interaction mediated by membrane-bound CYP3A4 (Denisov et al. 2015).

Determination of the Membrane-Binding Face of Hemoglobin N

Similar to cytochromes, hemoglobins are soluble proteins that are known to interact with the membrane (Fischer et al. 1975; Pathania et al. 2002). The primary function of hemoglobins is to bind small diatomic molecules, including O_2 and NO (Milani et al. 2001; Oliveira et al. 2012; Mairbäurl and Weber 2012), which, given their largely apolar nature, prefer to partition into the membrane (Subczynski and Hyde 1983; Shiva et al. 2001). Thus to maximize access to its binding partners, hemoglobins can take advantage of membrane interactions which permit substrate access through one of several pathways previously identified by Daigle et al. (2009). However, prior to a combined computational and experimental study by Rhéault et al. (2015), it was unknown what part of hemoglobin N from *Mycobacterium tuberculosis* interacted with the bilayer or which substrate access pathway would be preferred to permit efficient NO radical detoxification in a low O_2 environment as is found in lesions caused by tuberculosis.

Through the use of the MPEx tool (Snider et al. 2009) and PPM server (Lomize et al. 2012), Rhéault et al. (2015) generated initial potential membrane binding orientations for hemoglobin N above the membrane surface. Using the HMMM, Rhéault et al. (2015) were able to achieve

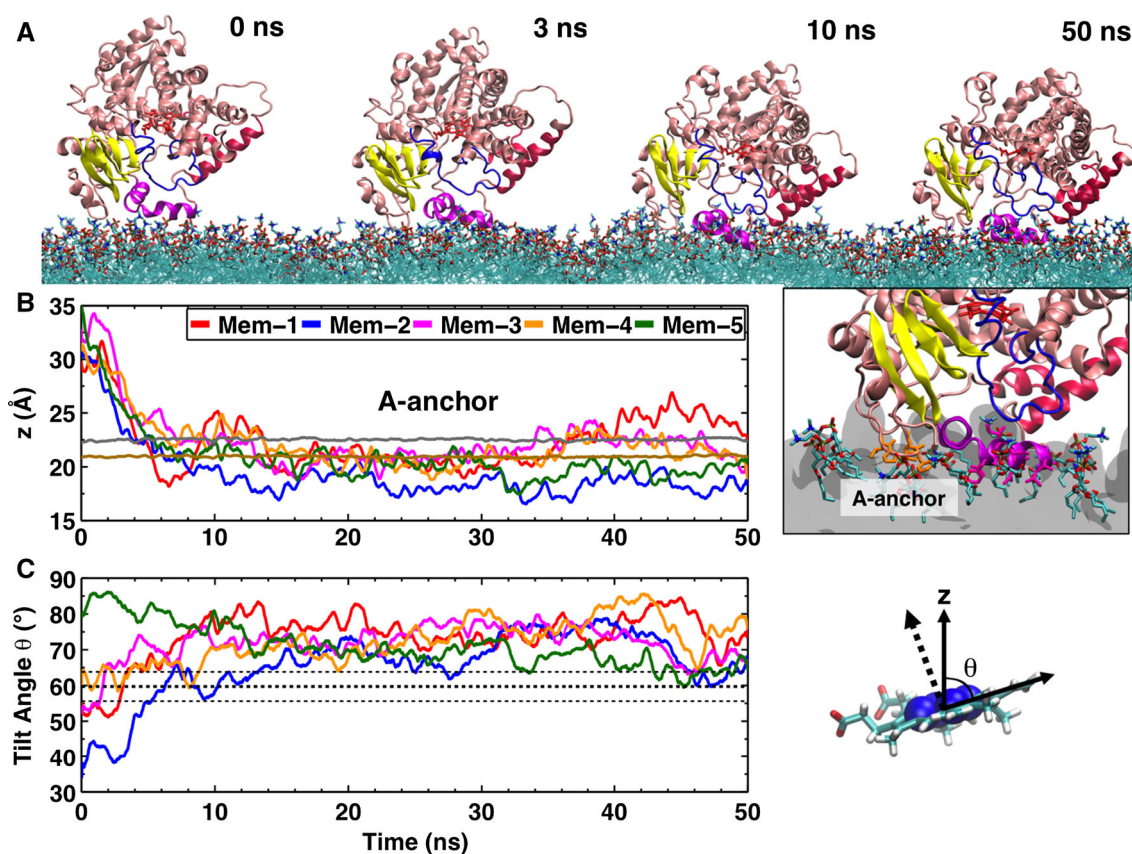


Fig. 7 Spontaneous membrane binding of CYP3A4 captured with the HMMM model. **a** Example trajectory of membrane binding. Membrane binding was consistently observed in five independent simulations (Mem-1 – Mem-5), permitting the identification of structural elements of CYP3A4 interacting with the membrane (e.g., A-anchor, F'-G' helices shown in magenta). **b** Insertion of the A-anchor into the membrane (i.e., at or below the PO₄ level) was consistently observed in all simulations, and converged to the same range of depths. The average positions of the phosphorus (PO₄) and the nitrogen (choline) atoms of the lipid head groups are shown using *brown* and *gray* lines,

respectively. **c** Each simulation began from a different orientation relative to the membrane, as measured by the heme tilt angle. The orientation for each of the simulations converged to around an average value of $72.2^\circ \pm 3.2^\circ$. The orientation of the protein was also measured experimentally, with an average tilt angle of $59.7^\circ \pm 4.1^\circ$ (*thick and thin dotted lines* for average and standard deviation in tilt angle plot, respectively), yielding excellent agreement between simulations and experiment. Adapted from Baylon et al. (2013). Copyright 2013 American Chemical Society (Color figure online)

converged binding depths and orientations, which did not change upon subsequent simulation in a conventional bilayer representation. This understanding meshed with experimental findings that the membrane surface was dehydrated by the approach of the protein as measured by IR spectroscopy, however, the simulations provided the atomic details of the interactions between the membrane and protein. This combined computational and experimental approach implicated the acidic residue Asp100 in modulating membrane binding via an electrostatic mechanism (Rhéault et al. 2015).

Membrane-Induced Structural Rearrangement of Talin

Integrins, as cell adhesion proteins, play an important role in many physiological and pathological processes,

including cell differentiation, platelet aggregation, and tumor metastasis (Hynes 2002; Liddington and Ginsberg 2002; Felding-Habermann et al. 2004; Ratnikov et al. 2005; Ma et al. 2007). These heterodimeric TM proteins are expressed in a default, low-affinity state that must be activated by the cytoskeletal-associated protein, talin (Tadokoro et al. 2003; Campbell and Ginsberg 2004; Ratnikov et al. 2005; Ma et al. 2007; Moser et al. 2009; Shattil et al. 2010). Talin is composed of two main domains: the rod domain, which interacts with the actin cytoskeleton, and the head domain, which interacts with the cell membrane and activates integrin (Critchley 2009). It is known that talin must bind to anionic phospholipids prior to activating integrin (Rees et al. 1990; Muguruma et al. 1995; Martel et al. 2001; Anthis et al. 2009; Saltel et al. 2009), demonstrating the importance of the membrane in the integrin activation process. While biochemical data

(Martel et al. 2001; Saltel et al. 2009) suggest that the talin F3 subdomain, the subdomain responsible for integrin activation, contacts the membrane surface, based on crystallographic structures, it was difficult to explain concomitant binding of both F2 and F3 subdomains to the membrane (Anthis et al. 2009; Elliott et al. 2010). Previous simulations (Kalli et al. 2010, 2013) utilized coarse-grained MD to describe the process of membrane binding of talin and might have not been able to fully describe some of the atomic details important to the process. Moreover, while biophysical studies had proposed the presence of a hydrophobic membrane anchor that inserts into the membrane (Dietrich et al. 1993; Tempel et al. 1995; Isenberg and Goldmann 1998; Seelig et al. 2000), the nature of such an anchor remained elusive.

To understand how the structural and dynamic changes in talin are modulated by membrane binding and to allow for optimal talin interaction with the membrane, Arcario and Tajkhorshid (2014) performed all-atom MD simulations with a PS HMMM membrane, resulting in a reproducible model of the membrane-bound talin F2F3 (Fig. 8). The starting talin structure was placed in five randomized initial orientations with varying heights above and angles with the PS HMMM bilayer. Across all five simulations, talin bound to the membrane in a consistent manner and in three distinct steps. First, a collection of basic residues in the F2 subdomain, known as the membrane-orientation patch (MOP), is attracted to the anionic membrane surface (Fig. 8b). While this patch had been identified as crucial for talin binding in previous mutagenesis experiments (Wegener et al. 2007; Saltel et al. 2009; Kalli et al. 2010), additional residues belonging to the MOP were identified via these simulations. Following initial association via the MOP, an initially buried, Phe-rich hydrophobic membrane anchor is exposed and inserted into the membrane core (Fig. 8c). Previous experimental studies predicted the existence of just such a hydrophobic anchor (Dietrich et al. 1993; Tempel et al. 1995; Isenberg and Goldmann 1998; Seelig et al. 2000), however, its identification and description had proven elusive both to structural and other computational studies (Anthis et al. 2009; Elliott et al. 2010; Kalli et al. 2010). Finally, a large-scale interdomain conformational change is induced by the membrane, leading to spontaneous binding of the F3 domain to the anionic membrane (Fig. 8d). This large conformational change produces a membrane-bound state in which the F2 and F3 subdomains interact with the membrane simultaneously, bridging the gap and reconciling previous discrepancies between biochemical and structural studies. Following these local and large-scale conformational changes, talin is optimally positioned to interact with integrin.

In order to confirm the validity of this membrane-bound model, Arcario and Tajkhorshid (2014) simulated the membrane-bound talin in a full-tailed, conventional

membrane. Starting from a membrane-bound configuration obtained from HMMM simulations, DOPS lipids were matched exactly to existing short-tailed PS HMMM headgroups and the tails oriented towards the center of the membrane, thereby preserving the specific lipid–protein interactions established in HMMM simulations. Over a series of successive steps, the membrane was relaxed around a fixed, membrane-bound talin. Following membrane relaxation, talin was simulated for 100 ns and remained stably bound throughout the simulation, retaining all the molecular interactions observed during the HMMM simulations.

Membrane Curvature Triggered by Synaptotagmin I

Neurotransmitter release is mediated by a Ca^{2+} -triggered membrane fusion event that fuses synaptic vesicles to the presynaptic membrane. During vesicle release, Syt1, acts as a Ca^{2+} sensor protein, binding to and bending the presynaptic membrane with its C2B domain, thereby regulating membrane fusion (Hui et al. 2009; Martens et al. 2007). With the help of MD simulations employing the HMMM model, Wu and Schulten (2014) proposed a molecular mechanism for how the C2B domain of Syt1 bends the presynaptic membrane and facilitates membrane fusion. Upon membrane binding, C2B moves its C-terminal helix towards the membrane, and has both its Ca^{2+} loop and the C-terminal helix shallowly inserted into an anionic membrane. This shallow insertion causes a pressure imbalance between the two membrane leaflets and thus causes membrane bending, which might facilitate the fusion event.

Wu and Schulten 2014 applied the HMMM model (Ohkubo et al. 2012) to the first stage of the simulations to accelerate C2B membrane binding and lipid redistribution around the domain, leading to spontaneous C2B binding within 30 ns. Importantly, this study is an example of a mixed bilayer HMMM application, where PS and PC headgroups were mixed in a 1:1 ratio to observe lipid redistribution of PS around the inserting C2B domain, which was found to concentrate around the Ca^{2+} -binding loops and lysine residues in the C-terminal helix. In simulations without anionic lipids, the Ca^{2+} loop does not bind to the membrane. Both the hydrophobic lipid tails and DCLE interact with the hydrophobic residues in the C-terminal helix, resulting in a $\sim 40 \times 15 \text{ \AA}^2$ region devoid of lipid headgroups around the helix in the cis-membrane leaflet (Fig. 9). Such lipid behavior has also been reported in many membrane-active peptides, such as melittin (Bernèche et al. 1998), influenza hemagglutinin (Lagüe et al. 2005; Larsson and Kasson 2013), and has been proposed to

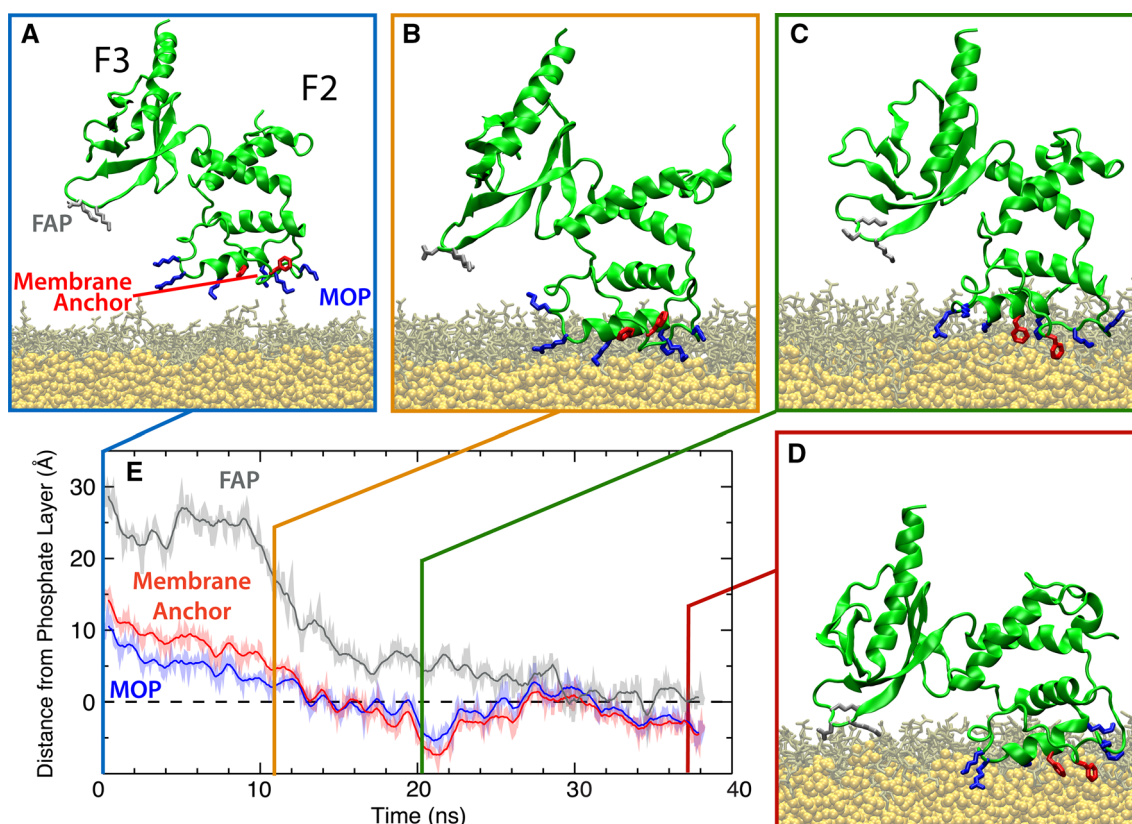


Fig. 8 Membrane-binding sequence of the talin F2F3 subdomain. **a** Initial configuration of the talin F2F3 subdomain (green) in the PS HMMM system (DVPS, brown DCLE, yellow). Residues comprising the MOP (blue), FAP (gray), and Phe-rich membrane anchor (red) are shown. In all images, water and ions have been omitted for clarity. **b–d** The three distinct stages of talin binding to the membrane are demonstrated. Talin is initially attracted to the membrane via MOP residues (**b**), followed by insertion of the Phe-rich membrane anchor

(**c**), and, finally, a large, interdomain conformational change that brings the F3 subdomain into contact with the membrane surface (**d**). **e** Timeseries of the height of the MOP, FAP, and Phe-rich membrane anchor with reference to the phosphate layer of the *cis* leaflet (dashed line) from a representative membrane-binding simulation are shown. See Arcario and Tajkhorshid (2014) for data on all five independent simulations (Color figure online)

facilitate membrane fusion (Smirnova et al. 2010; Risselada et al. 2011; Larsson and Kasson 2013).

In the second phase, the HMMM lipids were converted back to full lipids, in order to recover membrane mechanical properties necessary for reliable description of protein-induced membrane bending. Membrane bending within the HMMM is unusual in HMMM bilayers, as the organic solvents decouples the top and bottom leaflets from one another, and generally localizes bending to the *cis* leaflet. With the C2B domain already inserted in the first step, the C2B domain bends the presynaptic membrane spontaneously within 100 ns with a ~ 102 nm diameter membrane curvature. Due to the shallow insertion of the C-terminal helix into the membrane, only lipid head groups are displaced from the insertion site, while lipid tails in the *cis*-membrane leaflet approach the helix hydrophobic residues and become disordered. Coupling with the resulting void in the *cis* leaflet, lipids in the *trans* leaflet adopt a more stretched and ordered conformation. The difference

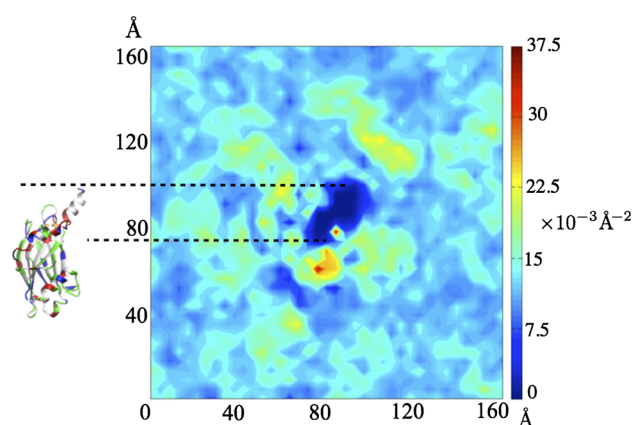


Fig. 9 Two-dimensional lipid head group density of the *cis* leaflet around synaptotagmin C2B domain (top view) resulting from simulations employing the HMMM model. The C2B domain with its protruding C-terminal helix is shown to scale on the left. The color of the density plot reflects the lipid occupancy per unit area, measured using the phosphorous atom of the lipid headgroup (Color figure online)

in lipid-tail ordering between the proximal and distal membrane leaflets generates a pressure imbalance across the bilayers, which serves as the driving force for C2B-driven membrane bending.

Membrane Origins of Structural Heterogeneity in α -Synuclein

α -Synuclein (α S) is a 14 kDa protein that is unstructured in solution, but forms helices when it binds to anionic membranes (Burré et al. 2013; Lokappa et al. 2014). Membrane-induced conformational changes of α S are of particular interest, given the connection between α S misfolding and aggregation and Parkinson's disease (Singleton et al. 2003; Charras et al. 2004). Experiments have identified two predominant states for monomeric α S bound to anionic membranes, a horseshoe-like broken helix identified with NMR and EPR (Ulmer et al. 2005; Drescher et al. 2008; Rao et al. 2010) where a kink brings together the two helical segments, and an extended kink-free helix conformation without a kink postulated by ESR and DEER experiments (Jao et al. 2004; Georgieva et al. 2008; Jao et al. 2008; Trexler and Rhoades 2009). It has been suggested that the extended and broken helix conformations could interconvert (Georgieva et al. 2010; Robotta et al. 2011; Lokappa and Ulmer 2011), potentially in a lipid-dependent fashion (Borbat et al. 2006; Middleton and Rhoades 2010). Outside of coarse-grained simulations (Braun et al. 2012; Braun and Sachs 2015), this conformational diversity had not been sampled, and presented a prime candidate for using the HMMM model to investigate the specific lipid interactions that determine the membrane-bound structure of α S.

Unlike in prior atomistic simulations where α S needed to be placed within the bilayer (Perlmutter et al. 2009), the HMMM model allowed Vermaas and Tajkhorshid (2014a) to observe spontaneous binding and insertion of α S to the bilayer and arrive at an unbiased pool of membrane-bound α S conformers. The pool had a mean insertion depth consistent with prior simulation and experimental findings (Jao et al. 2004; Perlmutter et al. 2009), although the variability observed in the HMMM systems was larger. Vermaas and Tajkhorshid (2014a) also noted that interaction between α S and PS broke apart interhelical hydrogen bonds between the N- and C-terminal helices in a subset of the simulations. Intercalation of PS between the helices provides alternative binding partners for basic residues, driving the helices apart (Fig. 10). In a conventional bilayer, mixing occurs too slowly to effectively sample the formation of PS microclusters underlying this process. In addition to using the HMMM to enhance lipid sampling, it was critical that the simulations were repeated several times, given the stochastic nature of the process. Only 20 % of the trials

demonstrated a semi-extended conformation where the inter-helix angle (θ_h) are greater than 47° (Fig. 10).

To distinguish these weak signals that occur infrequently, at least 10, and preferably 20 independent simulations need to be run to ensure that the observed behavior is a significant result. With the rapid insertion enabled by the HMMM, this level of statistical significance can be obtained with modest computational investment, estimated by Vermaas and Tajkhorshid (2014a) to be an order of magnitude less than what is required in a conventional lipid bilayer for the same degree of natural variability. Even with the advent of special purpose machines such as Anton (Shaw et al. 2008, 2014), using the HMMM allows researchers to see rare events at the membrane interface more frequently and gather superior statistics on the process of interest.

Acrobatics of the Synaptobrevin Transmembrane Domain

SNARE complexes, the machinery of neuronal exocytosis, are composed of fusion proteins located at both the plasma (syntaxin, SNAP-25) and vesicular (synaptobrevin) membrane surfaces (Fasshauer et al. 1997; Südhof 2004; Südhof and Rothman 2009). These proteins initially embedded in different (vesicular and plasma) membranes must intimately interact, forming a 120 Å-long triple helix, in order to prime the neurotransmitter-filled vesicle for fusion with the neuronal plasma membrane (Sutton et al. 1998). Synaptobrevin comprises the vesicular component of the SNARE complex and is responsible for targeting a particular vesicle to the plasma membrane and priming it for membrane fusion. While the post-fusion structure of the SNARE complex is well studied (Brunger 2005; Jahn and Scheller, 2006; Südhof and Rothman 2009; Stein et al. 2009) and known to stand parallel with the membrane normal (Sutton et al. 1998) with multiple interhelical interactions apparent (Brunger et al. 2009), the details of the pre-fusion state of the individual SNARE proteins is not as well understood.

Studies on the transmembrane (TM; residues 95–116) and linker (residues 85–94) regions have suggested a large tilt ($\sim 30^\circ$ – 50°) of the linker region of synaptobrevin with respect to the membrane normal (Kweon et al. 2003; Bowen and Brunger 2006). Conflicting theories on how to explain this large tilt have emerged, however, with some expressing that the linker region is a continuous helix (Bowen and Brunger 2006) and some arguing for a disordered structure (Durrieu et al. 2009) allowing flexibility between the TM and cytosolic (residues 1–84) domains. The structure of the TM and linker regions is extremely important to the function of synaptobrevin as minor alterations drastically affect the fusion potential of the SNARE

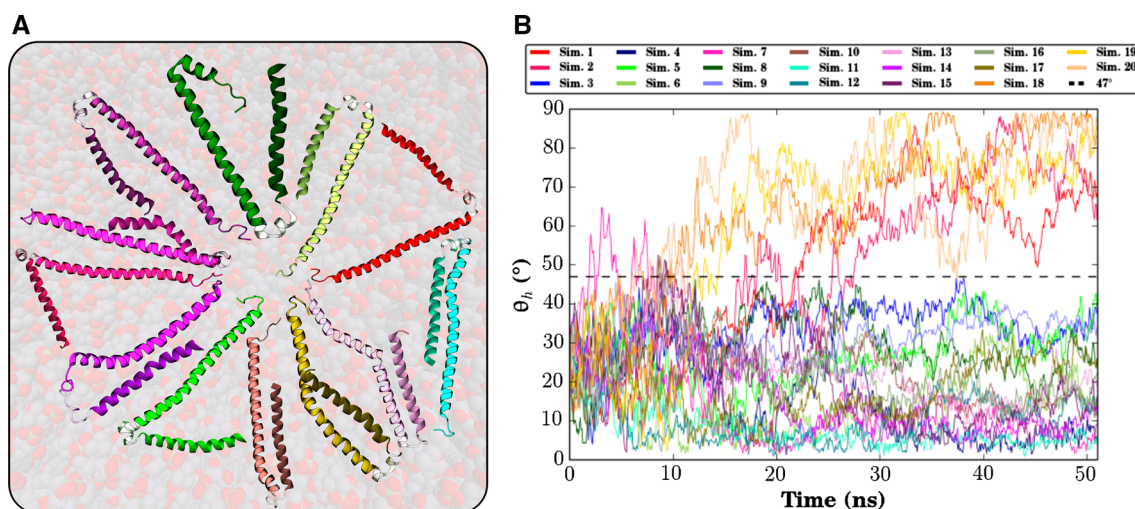


Fig. 10 **a** Selected end states of the α S monomers across the 20 HMMM membrane-binding simulations conducted, highlighting the variability of α S when bound to the membrane. **b** Inter-helix angle (θ_h) for all 20 membrane-binding simulations. θ_h is defined by the angle between the two vectors formed by a linear interpolation of the

α carbons of the N- and C-helices. θ_h distinguishes between two states, a broken helix state where θ_h never exceeds 47° , and a semi-extended intermediate state where the inter-helix angle is approximately 90° . Originally published in Vermaas and Tajkhorshid (2014a), used with permission from Elsevier B.V

complex (McNew et al. 1999; Kweon et al. 2003; Ngatchou et al. 2010). The challenge in finding the spontaneous native tilt of synaptobrevin using MD simulations, however, is the extended time scale associated with lipid diffusion. Tilting of a TM helix requires significant rearrangement of lipids in both membrane leaflets.

To overcome these limitations, Blanchard et al. (2014) carried out simulations using the HMMM model with short-tailed PC lipids to attempt to distinguish between these two hypotheses. This study represents the first example of a TM application of the HMMM model. Five independent simulations, with the initial tilt angle of synaptobrevin ranging from -20° to $+40^\circ$, were performed, each simulation converging to a tilt angle of $40.1 \pm 3.4^\circ$ within 50 ns (Fig. 11b). The tilt angle resulting from these simulations is within the range provided by previous structural studies (Kweon et al. 2003; Bowen and Brunger. 2006), and the fact that the independent simulations converge from a wide range of starting configurations removes the possibility that the final tilt angle was biased by its initial orientation. Blanchard et al. (2014) also identified specific motifs of charged and aromatic residues that dictate the native tilt of synaptobrevin, offering mechanistic insight into why previous mutants had diminished fusion ability (Kweon et al. 2003).

Additionally, because of the fluidity of the DCLE layer, Blanchard et al. (2014) were able to observe significant dynamics of the TM segment of synaptobrevin on the MD timescales that contain structural implications for previous experimental findings. In particular, while the “linker” region remains completely helical, consistent with circular dichroism (CD) studies (Bowen and Brunger 2006), a

glycine residue at the 100 position disrupts the helical structure and allows the lower portion of the TM domain (TM2) to move independently of the upper portion (TM1) (Fig. 11a). Thus, while TM1 adopts a clearly defined angle relative to the membrane normal, TM2 does not and explores much more space. This exact behavior had been proposed by Bowen and Brunger (2006) as a possibility to explain large discrepancies in CD signals between the “linker” and TM regions of synaptobrevin. The structural findings made possible by the HMMM remained unchanged when Blanchard et al (2014) built a conventional bilayer model starting from the HMMM endpoint, highlighting the general applicability of the HMMM to study particular elements of TM systems.

Best Practices

If these example applications have inspired you to use the HMMM in your own research, there are several key points where the setup of HMMM simulations is different from a conventional membrane simulation. The most consequential difference is that use of the HMMM requires a membrane with fixed surface area, either by running at constant volume or a constant area barostat. In a conventional constant ratio barostat, the surfactant-like headgroups may cause membrane expansion (Fig. 12a). For membrane-binding studies, this expansion is beneficial to a degree, as it both increases lateral lipid diffusion modestly (Fig. 12b) and creates space for an incoming protein. However, if left unchecked the expansion may result in a sparse membrane, and so prior HMMM simulations used a fixed area

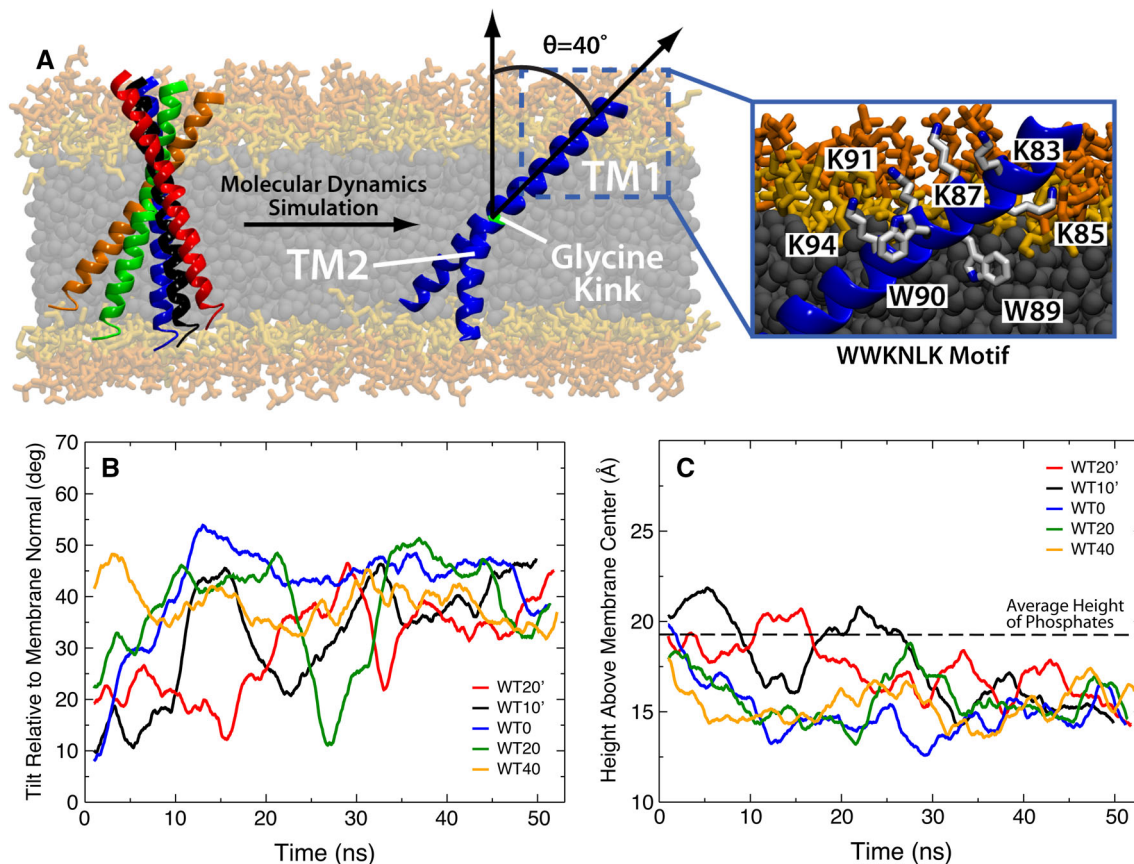


Fig. 11 Native tilt angle of the synaptobrevin TM domain. **a** The five initial conditions of synaptobrevin tilt simulated are shown together with the final structure. All simulations converged to an average tilt of $40.1 \pm 3.4^\circ$ with the two segments of the helix (TM1 and TM2) mechanically decoupled by the G100 “glycine kink”. A detailed image of the specific interactions leading to the well-defined orientation of TM1 are shown in the *inset*. Plots of the angle relative

to membrane normal (**b**) and height of tryptophan residues above membrane center (**c**) are shown for each of the five simulations, with the color of the trace in *each plot* matching the color of the helix in (**a**). In **c**, the *black dashed line* represents the average height of phosphates in the leaflet closest to the tryptophans (upper leaflet in **a**) across all simulations

ensemble, with the pressure controlled only by varying the dimension parallel to the membrane normal (Ohkubo et al. 2012; Baylon et al. 2013; Vermaas and Tajkhorshid 2014a, 2014b; Blanchard et al. 2014; Arcario and Tajkhorshid 2014). A good starting area per lipid considers the estimated volume of lipids displaced by the incoming protein. Even though it is not possible to know a priori how deep the penetration of the protein into the membrane will be, a reasonable rule of thumb is that the area per lipid initially should be approximately 5–10 % higher than what is observed in a conventional bilayer, which results in a final effective area per lipid that is nearly physiological (Baylon et al. 2013; Vermaas and Tajkhorshid 2014a).

In dense membranes, the short acyl tails of the head-group surfactants may also cause the headgroups to go into solution. While partitioning into the aqueous phase is a natural surfactant-like behavior expected for short-tailed lipids, this is generally undesirable for a bilayer mimic, as it reduces the number of lipids in the membrane. In order to

address this behavior, a mild $0.1 \text{ kcal/mol/\AA}^2$ restraint along the membrane normal can be applied to the carbonyl carbon of the acyl tails (C21 and C31 in CHARMM parlance). With this restraint, the z-positions of the membrane components still readily fluctuate by approximately 3.5 \AA in each direction at room temperature, resulting in excellent overlap between the membrane profiles of corresponding atoms in HMMM and conventional bilayers. As for the requirement for fixed area ensembles, these recommendations can be adjusted to suit the system, such as using a physiological lipid density for TM peptides which are already inserted (Blanchard et al. 2014).

For larger TM systems than those discussed in this review, special care needs to be taken to avoid intercalation of organic solvent into the TM domain. Since DCLE is slightly polar, it can and will disrupt the interactions that hold together TM proteins if left unchecked. This can be addressed by specialized *in silico* solvent models currently in development. However, for many applications it is

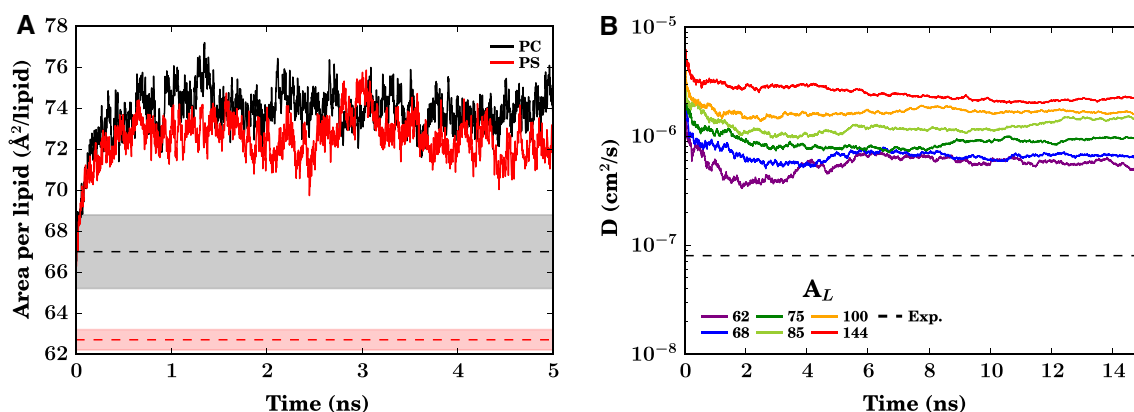


Fig. 12 **a** The evolution of the area per lipid of HMMM bilayers in a constant ratio barostat for both PC (black) and PS (red) headgroups, highlighting the bilayer expansion possible due to the shorter acyl tails. For reference, the range of experimental areas per lipid (A_L) determined through a combination of X-ray and neutron scattering methods by Fogarty et al. (2015) are reported as shaded regions, with the mean as a dashed line. **b** Lateral lipid diffusion constant (D_L) time

evolution for PC HMMM bilayers over a range of A_L maintained by a fixed-area barostat. The dashed line in this panel represents the D_L of $\sim 8 \times 10^{-8} \text{ cm}^2 \text{ s}^{-1}$ for DPPC lipids (Klauda et al. 2006; Wohlert and Edholm 2006), consistent with most conventional bilayers. The simulations behind this data were conducted using NAMD 2.10 (Phillips et al. 2005) on bilayers constructed using CHARMM-GUI (Jo et al. 2008, 2009; Wu et al. 2014) at 310 K (Color figure online)

sufficient to restrain the TM domain to retain its shape through the use of harmonic restraints on the atomic positions to keep the protein packed correctly. Such restraints have minimal bearing on the lateral diffusion rate of the lipids around the protein, and do not prevent specific lipid-binding sites around the protein from being explored.

In terms of constructing HMMM systems, the step by step procedure is virtually identical to how one would construct a conventional bilayer. For typical situations, where the interest is on sampling an unbiased range of conformations, it is important to generate a diverse starting pool of initial bilayers so that the final state is not unduly biased by the starting geometry of the bilayer. The simplest means of manufacturing HMMM bilayers with diverse lipid compositions is to use the HMMM Builder as recently implemented in CHARMM-GUI (Jo et al. 2008, 2009; Wu et al. 2014). An HMMM bilayer can also be constructed from a conventional bilayer by replacing the entire acyl tail beyond carbon five with an organic solvent. Special care, however, should be taken in packing in the right amount of organic solvent, as overfilling or underfilling the bilayer interior can lead to unstable bilayers. In order to match membrane volumes between HMMM and conventional bilayers, 3 DCLE molecules should be added for every 14 acyl tail carbons that have been removed, plus any additional adjustments that need to be made to create a higher A_L if desired

$$M = \frac{3}{14} C \frac{A_L^{\text{HMMM}}}{A_L^{\text{conventional}}}$$

Here, M is the number of DCLE molecules to be added to the bilayer center, C is the number of acyl tail carbons removed from the conventional bilayer, and $\frac{A_L^{\text{HMMM}}}{A_L^{\text{conventional}}}$ is the

ratio of the A_L (area per lipid) between the target HMMM bilayer and the conventional bilayer initially generated.

The final point is that the HMMM excels in the initial sampling phases, but should be converted back into a conventional bilayer at the end prior to extensive equilibrium simulations. In general, it is preferred that the existing short acyl tails from the HMMM lipids are preserved during this conversion process, since the existing tails will not already clash with the other membrane components. Vermaas and Tajkhorshid (2014a) describes an additional procedural enhancement through the use of the organic solvent heavy atom coordinates to seed the positions of the acyl tails around the inserted species; however, in many instances, using internal coordinate tables to extend the tails results in acceptable starting geometries for continued simulation (Baylon et al. 2013; Wu and Schulten 2014).

Future Directions

Along with experimental characterization of additional peripheral proteins whose binding to membranes is controlled by the lipid composition and with functional assays further establishing the regulatory roles of specific lipid molecules in membrane protein function in general, the HMMM has a bright future as a method to complement these findings with a mechanistic description of the specific lipid-protein interactions at work. The model itself is also expected to evolve, both through the addition of designer *in silico* solvents that ameliorate deficiencies in the representation of the hydrophobic cores of the membrane, and new problems in membrane biology to which the HMMM could be applied.

Many of these applications move beyond the realm of peripheral proteins, e.g., into TM systems. One such application is to use the HMMM to identify lipid-specific binding sites on integral membrane proteins with affordable computational effort, as the accelerated lateral lipid diffusion would allow for more rapid sampling on easily accessible computational timescales. Other pressing challenges in membrane biophysics, such as ion-induced lipid clustering and protein-induced enrichment of specific lipids in its surrounding micro-environment, would continue to significantly benefit from the rapid sampling of lipid-protein interactions the HMMM provides. In particular, when used in conjunction with the recent advent of routine μ s simulation, the HMMM can begin to tackle challenges long beyond the reach of conventional MD, such as the assembly of membrane-embedded super-complexes at an atomic level of detail. Thus, the number of applications of the HMMM will only continue to grow into the future.

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